### 2. Methods

### 2.1 Bubble liposomes

We used Bubble liposomes prepared with polyethyleneglycol-modified liposomes (PEG-liposomes) and perfluoropropane gas (Takachiho Chemical Inc., Co., Ltd., Tokyo, Japan). The liposomes were prepared by a reverse-phase evaporation method, as described previously[10]. To prepare liposomes for the BLs, 1, 2-distearoylsn-glycero-phosphatidylcholine (DSPC) and 1, 2distearoyl-sn-glycero-3-phosphatidyl-ethanolamine-methoxypolyethyleneglycol [DSPE-PEG (2k) -OMe] were mixed at a molar ratio of 94:6. First, 5 mL sterilized vials containing 2 mL of PEG-liposome suspension (lipid concentration: 1 mg/mL) were placed in vials supercharged with 7.5 mL of perfluoropropane (C<sub>3</sub>F<sub>8</sub>) gas, then sonicated by continuous ultrasound at a frequency of 42 kHz for two or three min. During this process, gas was trapped inside the liposomes, resulting in a cloudy suspension. A bath-type ultrasound cleaner (Branson 2510) was used for sonication. Figure 1 shows an image of BLs under a microscope equipped with a Darklite Illuminator (NEPA Gene Co. Ltd., Chiba, Japan). Figure 2 shows the size distribution of BLs and Sonazoid measured by dynamic light scattering (ELS-Z, Otsuka Electronics Co., Ltd., Osaka, Japan). Although sub-micron sized objects are not shown in the microscopic image, Fig. 2 indicates that the average diameters of BLs are approximately 400 to 500 nm and are smaller than the conventional MBs Sonazoid. Table 1 shows the characteristics of BLs and another type of MBs F-04E[6, 7]. The suspension of BLs was freshly prepared before the experiment,

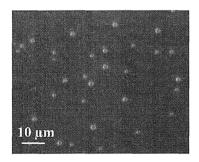


Fig. 1 Microscopic image of BLs.

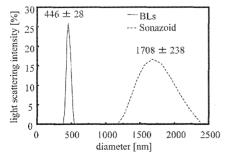


Fig. 2 Size distribution of BLs and Sonazoid.

diluted with saline to a concentration of  $0.01\text{--}0.05\,\mathrm{mg}$  lipid/ml.

## 2.2 Experimental setup

We prepared an artificial blood vessel with a Y-form bifurcation structure, which was made of a mixture of wax and poly (vinyl alcohol) (PVA) [11]. The inflow path of 2 mm was repeatedly divided into two lower courses to provide artificial capillaries until the middle of the model. where the minimum path width was 0.50 mm. The path widths and the cross-sectional areas were designed to guarantee a constant flow velocity in any part of the model. The whole view, x-y plane view and x-z plane view of the experimental setup are shown in Fig. 3a, b and c. respectively. The artificial blood vessel with external dimensions of  $180 \times 70 \times 8 \text{ mm}^3$ , was positioned 30 mm above the bottom of a water tank, to prevent multiple reflections of ultrasound between the artificial blood vessel and the bottom of the tank. The blood vessel was divided into Paths A and B with a branch angle of 75 degrees. The path widths  $w_1$  and  $w_2$  corresponded to 1.4 mm and 1.0 mm, respectively. We used an optical microscope (Hirox KH-7700) to observe the area of interest in the blood vessel. The image size was  $800 \times 600$ pixels and the optical resolution of the digitized images was  $31.5 \,\mu m$  per pixel. The light source was located above

Table 1 Characteristics of BLs and MBs.

	BLs	MBs(F-04E[6, 7])
Mean diameter $[\mu \mathrm{m}]$	0.4-0.5	2.7
Shell	Lipid(bilayer)	PVC-AN copolymer
Gas	$C_3F_8$	$C_3H_8$ , $C_4H_{10}$

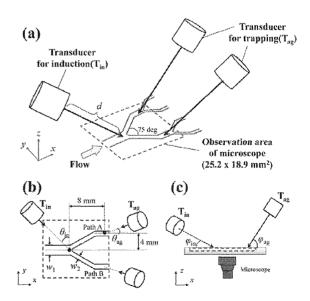


Fig. 3 Schematic presentation of the experiment with the artificial blood vessel, the microscope and the ultrasound transducers.

the water tank. We set transducer T<sub>in</sub>, consisting of a concave ceramic disc, to emit focused ultrasound for induction. We prepared four different ultrasound transducers Tin to compare the effect of frequency. The central frequencies of the T<sub>in</sub> were 3, 5, 7 and 10 MHz, and the apertures were 18, 15, 12 and 10 mm, respectively. The focal lengths of individual transducers ranged from 53 to 62 mm. Kudo et al. [12] reported stability of oscillating BLs under wideband ultrasound pulses at a central frequency of 10 MHz. We investigated frequencies below 10 MHz, considering the frequency range of clinical ultrasound. The focal position of  $T_{\rm in}$  was fixed at the bifurcation point. The angle of the axis of  $T_{in}$  was set at  $\theta_{in} = 45$  degree and  $\varphi_{in}$  = 30 degree. In our previous study [7], we investigated the controllability of MBs in a sound pressure range of 100 to 500 kPa-pp because of output bound of the elements of the transducers. In the present study, we set the sound pressures at 100, 200, 300, 400 and 500 kPa-pp.

Other transducers Tag were newly introduced for evaluation of the controllability of BLs, as mentioned in the next section. The transducer Tag was the same as one of the T<sub>in</sub>, which consisted of a concave ceramic disc with a central frequency of 5 MHz and an aperture of 15 mm. The focal length of the Tag was 58 mm. We set the axis of one  $T_{\text{ag}}$  to correspond to the center of the axis of the Path A. The focal point of  $T_{ag}$  was located 8 mm on the x-axis and 4 mm on the y-axis from the bifurcation point. The other T<sub>ag</sub> was located in the middle of Path B (8 mm on the x-axis and 4 mm on the y-axis from the bifurcation point). We set the maximum sound pressure of T<sub>ag</sub> at 500 kPa-pp. The angle of the axis of  $T_{ag}$  was set at  $\theta_{ag} = 20$ degree and  $\varphi_{ag}$  = 60 degree. The distance of the transducers from the blood vessel was set at d = 60 mm. Sinusoidal waves with frequencies of 3, 5, 7 and 10 MHz were generated by an oscillator and applied to transducers Tin and Tag. Figure 4 shows an overview of the experimental setup.

# 2.3 Measurement of the area of trapped BLs

In our previous studies [6, 7, 11], we evaluated the controllability of MBs by defining the induction index that reflects the brightness according to the presence of

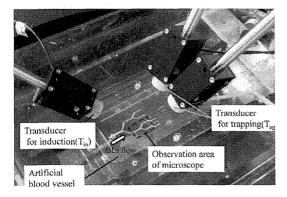


Fig. 4 Overview of the experimental setup with transducers and the artificial blood vessel.

bubbles in two paths (Path A and B). However, because of the size and low concentration of BLs, observation of brightness was difficult. Thus, we calculated the area of trapped BLs in each path and defined a new induction index as follows. To measure the number of BLs, we established three square regions of interest (ROI) in the path before bifurcating (ROI O) and in Paths A and B (ROIs A and B), respectively. Figure 5 shows the configuration between transducers T<sub>ag</sub> and the ROIs. The transducer drawn in dotted line (Tag'), which was the same type of transducer as the other two Tag, targeted a point 5 mm upstream from the bifurcation point. The angles of the axis of  $T_{ag}$  were set at  $(\theta_{ag}, \varphi_{ag}) = (0, 60)^{\circ}$ . This configuration was set to confirm the degree of conservation of the number of BLs before and after the hifurcation.

The width of each ROI corresponded to the path width, and the length of each ROI (L) was set by considering the ultrasound beam width. In our previous research [8, 9], the half-widths of ultrasound beams at 3 and 10 MHz were 1.7 and 1.1 mm, respectively. We set L and  $L_0$  at 6.5 and 10.0 mm, which were much larger than the ultrasound beam width. Then, when the size of ROI A or B was  $1.1 \times 6.5 \, \mathrm{mm}^2$  and the size of ROI O was  $1.6 \times$ 10.0 mm<sup>2</sup>, we calculated the area of trapped BLs. To evaluate the controllability of the amount of BLs quantitatively, image processing[8] was used to calculate the area of BL aggregates trapped by Tag in the middle of the path. Figure 6 shows the image processing procedure for measuring the areas of trapped BLs in ROIs A and B. The outline of the blood vessel was overlapped as the dotted lines in the images. The microscopic image was recorded continuously as a video file from the beginning of the experiment. Then the area of trapped BLs was extracted by subtracting the initial image from the images after aggregates of BLs were formed. Finally, the area of trapped BLs was obtained from the binary image by discriminant analysis method.

**Figure 7** shows the changes over time of the areas of trapped BLs in ROIs A and B upon injection of the BL suspension with ultrasound emissions from both  $T_{\rm in}$  and  $T_{\rm ag}$ . The maximum sound pressure of  $T_{\rm in}$  was 200 kPa-pp and the concentration of BLs suspension was 0.02 mg lipid/ml. Before injection, the trapped area remained at

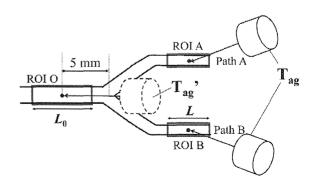


Fig. 5 Configuration between transducers T<sub>ag</sub> and ROIs.

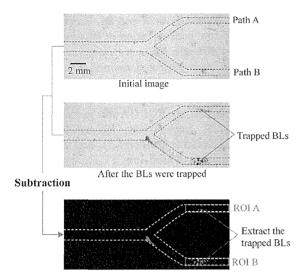


Fig. 6 Procedure to measure the area of trapped BLs in ROIs A and B.

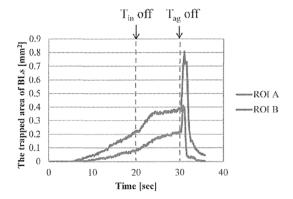


Fig. 7 Changes of the areas of trapped BLs in ROI A and ROI B over time.

zero. With the appearance of BLs, aggregates of BLs were trapped in both ROIs and the trapped areas increased simultaneously. A significant difference between the two ROIs was confirmed with T<sub>in</sub> emission, indicating that larger amounts of BLs were propelled by Tin ultrasound emission. However, no significant difference between the two paths was observed without Tin emission (data not shown). A 3-ml aliquot of BL suspension was injected into the flow. The duration of injection was controlled by a rotary pump, set to finish approximately 20 s after the beginning of injection at a flow velocity of 40 mm/s. Then T<sub>in</sub> emission was stopped at 20 s. Actually, all injected BLs were observed to have been delivered during 20-25 s. After 20 s, the trapped areas increased slightly, because the aggregates of BLs trapped at the bifurcation by T<sub>in</sub> were removed and trapped at ROIs A and B. At 30 s. emissions by Tag were stopped simultaneously, and thereafter the trapped area reached a maximum because the aggregates of BLs trapped by Tag collapsed and apparently dispersed. At 30-40 s after the beginning of measurement, the trapped areas declined due to dis-

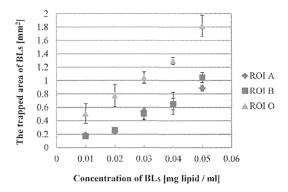


Fig. 8 The trapped area of BLs vs. concentration of BLs.

appearance of bubbles. Thus, we evaluated the area of trapped BLs at 20 s and 30 s, just before  $T_{\rm in}$  and  $T_{\rm ag}$  emissions were stopped, respectively. At 20 s and 30 s, we calculated the induction index  $\xi_B$ , which indicates that BLs are induced to Path B rather than Path A, using the following equation:

$$\xi_B = \frac{\sigma_B - \sigma_A}{\sigma_A + \sigma_B} \tag{1}$$

where  $\sigma_A$  and  $\sigma_B$  are the areas of trapped BLs in ROIs A and B, respectively.

Then, we confirmed the relationship between trapped area  $\sigma$  and concentration of BLs. To confirm the conservation of the number of BLs before and after the bifurcation, we measured the trapped areas in ROI O, ROI A and ROI B without ultrasound emission from Tin. Suspensions of BLs were prepared at five levels of concentrations ranging from 0.01-0.05 mg lipid/ml, and trapped in each ROI with the emission of T<sub>ag</sub> and T<sub>ag</sub> as shown in Fig. 5. The areas of trapped BLs were measured for suspensions of various lipid concentrations with ultrasound emission at the maximum sound pressure of 500 kPa-pp. **Figure 8** shows the concentration of BLs vs the area of trapped BLs. The standard deviation of three trials is also shown in the figure. The trapped areas in ROI A and B were measured simultaneously, the trapped area of ROI O was measured independently.

As shown in **Fig. 8**, using BL concentrations of 0.01-0.05 mg lipid/ml, the area of BLs increased linearly as the concentration of BLs increased. The trapped area in ROI O was approximately two-fold larger than that in ROI A or B. We confirmed that the concentration of 0.02 mg lipid/ml, which was used in subsequent experiment, was within the range of linear change in trapped area.

# 3. Results

We observed BL behavior in the bifurcation upon emission of ultrasound from  $T_{\rm in}$  and  $T_{\rm ag}$ . Figure 9 shows the serial extraction images of BLs in the observation area after the appearance of injected BLs with  $T_{\rm in}$  emission at a frequency of 5 MHz, a flow velocity of 40 mm/s, and at a BL concentration of 0.02 mg lipid/ml. The outline of the blood vessel was overlapped as the white line in the top-left image. The bottom-left image is a close-up image of bifurcation area enclosed by the dotted line in the image

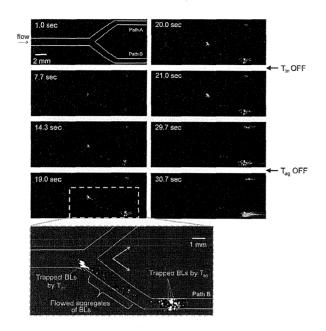


Fig. 9 Serial extraction images of BLs in the observation area.

of 19.0 s. We confirmed flowed aggregates of BLs in Path B after Tin emission although they were not observed before the emission. The aggregated BLs were trapped both in ROIs A and B. Up to 20 s, an increase in aggregation in ROIB was observed. On the other hand, an increase in area of BLs at the focal point of Tin in the bifurcation was also confirmed. When the ultrasound from T<sub>in</sub> was stopped, most the BL aggregates trapped by T<sub>in</sub> at the bifurcation collapsed and flowed into ROIs A and B. Most of the aggregates were trapped in the two ROIs in the duration from 20 s to 30 s, with additional increases in area of trapped BLs in both ROIs. After 30 s, the aggregates of BLs collapsed and disappeared.

We measured the areas of trapped BLs in ROIs A and B upon T<sub>in</sub> emission at a frequency of 5 MHz, for a BL concentration of 0.02 mg lipid/ml, when the sound pressure was set at 100, 200, 300, 400 and 500 kPa-pp. Figure 10 shows the areas of trapped BLs versus sound pressure of  $T_{in}$  at 20 s(a), and at 30 s(b). The standard deviation of three trials is also shown in the figure. Statistical significance was analyzed by t-test (\*P < 0.01, \*\*P < 0.05, \*\*\*P < 0.1). At 20 s, the area of trapped BLs in Path B  $(\sigma_B)$  remained unchanged at  $0.2 \text{ mm}^2$  when the sound pressure was 300 kPa-pp or higher (Fig. 10a). because more BLs were trapped at the focal point of T<sub>in</sub> at the bifurcation. At 30 s, the  $\sigma_B$  increased rapidly with an increase in sound pressure of T<sub>in</sub> (Fig. 10b). Comparing **Fig. 10a and 10b**, the  $\sigma_B$  increased 1.35, 1.87, 1.81, 1.88 and 2.03 times at sound pressures of 100, 200, 300, 400 and 500 kPa-pp, respectively. When the sound pressure exceeded 300 kPa-pp, the induction indexes at 30 s were lower than those at 20s.

Next, we fixed the sound pressure at 200 kPa-pp, and repeated the same experiment with the central frequencies at 3, 5, 7 and 10 MHz. Figure 11 shows the area of

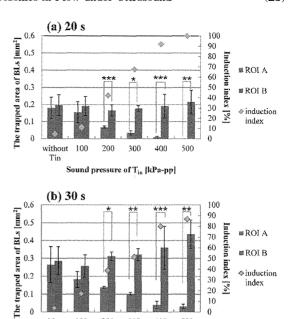


Fig. 10 The trapped area of BLs vs. sound pressure of  $T_{in}$  at 20 s (a), and at 30 s (b) with induction index. (P < 0.01, P < 0.05, P < 0.1).

300 400

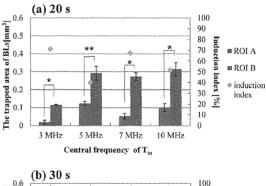
Sound pressure of T<sub>in</sub> [kPa-pp]

without 100 200

Tin

10

500



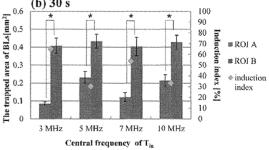


Fig. 11 The trapped area of BLs vs. central frequency of Tio at 20 s (a), and at 30 s (b) with induction index. (\*P < 0.01, \*\*P < 0.05).

trapped BLs versus central frequency. The standard deviation of three trials is also shown in the figure. Statistical significance was analyzed by t-test (\*P < 0.01, \*\*P < 0.05). Comparing Fig. 11a and 11b, the  $\sigma_B$ increased 3.52, 1.49, 1.48 and 1.36 times at central frequency of 3, 5, 7 and 10 MHz, respectively. At 20 s,

especially at the frequency of  $3\,\mathrm{MHz}$ , more BLs were trapped at the focal point of  $T_{in}$  in the bifurcation, with small trapped area in ROI B. Induction indexes at  $30~\mathrm{s}$  were all lower than those at  $20\mathrm{s}$ .

#### 4. Discussion

From the result of **Fig. 9**, when the T<sub>ag</sub> emission stopped, the aggregates of BLs in ROIs A and B collapsed. In our previous study [9], the aggregates of MBs Sonazoid maintained the shape even if after the ultrasound emission stopped, and therefore we can apply Sonazoid to a novel therapy, artificial embolization. This is a very different aspect between BLs and conventional MBs. In other words, there is no concern over embolization in ultrasound therapy using BLs. This difference seems to arise from the shell properties of BLs and Sonazoid. Sonazoid bubbles are bound to each other because of surface interaction of the lipid membrane without PEG modification, resulting in long-lasting bonding with each other.

BLs with average diameters of 400 to 500 nm should be subjected to 16 to 25 times less force than MBs that had an average diameter of  $2\,\mu\text{m}$ , because the primary Bjerknes force [7, 8] is proportional to the cross-sectional area of the objects. However, we succeeded to trap BLs at the same range of sound pressure as for MBs in this experiment. It was probably the effect of aggregation caused by secondary Bjerknes force under ultrasound emission. Because of the increased cross-sectional area caused by aggregation, the aggregates of BLs would be propelled easily. In terms of trapping efficiency, we have to consider not only the differences in diameter and shell properties, but also aggregate formation and interaction between surface of bubbles and vessel walls.

The induction index was proportional to sound pressure (Fig. 10). This tendency was consistent with MBs[6, 7]. Further study is required to clarify the sound pressure level for destruction of BLs. Under the threshold of sound pressure for shell destruction, BLs may be used as a carrier for gene and drug delivery until they arrive at their final destination inside the body. Furthermore, the threshold is very important for drug release.

More BLs were trapped on the vessel wall at the bifurcation point by exposure to  $T_{\rm in}$ , especially at the frequency of 3 MHz (**Fig. 11**). This occurred probably because of the wider beam width of ultrasound. The overlapped area of ultrasound beam on the vessel wall becomes larger as frequency becomes lower. Furthermore, the size of aggregates increases at lower frequencies [13], apparently increasing the amount of trapped BLs. To avoid trapping BLs at the bifurcation, thinner beam width and higher frequency are required.

In the next step, we will elucidate the behavior of BLs in a viscous flow before applying it to an *in vivo* experiment. The effect of hematocrit of human blood should be considered to determine the viscosity of fluid medium. We will also develop a method of identifying fluorochrome-labeled BLs *in vivo*.

## 5. Conclusions

In this study, we achieved active control of BLs with diameters of 400 to 500 nm, orientating these BLs to flow into the desired path of a bifurcated blood vessel. To measure the controllability of BLs quantitatively, we designed a new method using two transducers to evaluate the areas of trapped BLs in the two paths after the bifurcation, to determine which path has increased BLs. We investigated the ultrasound conditions for active path selection of BLs, in terms of the sound pressure and frequency. We found that more BLs can be oriented to a desired path at higher sound pressure. To avoid trapping BLs at the bifurcation, higher frequency is required. For further analysis, we are aiming at active control of BLs in vivo.

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### Ren Koda

Ren Koda received the B.Eng. and M.Eng. degrees from Chiba University, Japan in 2009 and in 2011, respectively and the Ph.D. degree in Engineering from Tokyo University of Agriculture and Technology, Japan in 2014. He is currently a Research Assistant



Professor of Science and Engineering for Research at University of Toyama, Japan. His research interests include ultrasonic therapy using microbubbles and non-destructive inspection using ultrasound.

# Jun Koido

Jun Koido received the B.E. degree in Engineering from Tokyo University of Agriculture and Technology, Japan in 2013. His research interest is microbubble delivery using ultrasound in blood vessel for his master degree.



## Naoto Hosaka

Naoto Hosaka received the B.Eng. and M.Eng. degrees from Tokyo University of Agriculture and Technology, Japan in 2013 and in 2014, respectively. He is currently a Doctorial course student, Tokyo University of Agriculture and Technology, Japan. His



research interest is microbubble manipulation in blood vessel for therapy.

#### Shinya Onogi

Shinya ONOGI received the B.Eng. and M.Eng. degrees from Hokkaido University, Japan in 2002 and in 2004, respectively, and the Ph.D. degree in Science from the University of Tokyo, Japan in 2007. From 2007 to 2009, he was a Project Researcher of



Intelligent Modeling Laboratory at the University of Tokyo. From 2009 to 2010, he was a Postdoctoral Fellow of Mechanical Engineering at Johns Hopkins University, USA. He is currently a Specially Appointed Assistant Professor (NEXT Program, JSPS) of Bio-Applications and Systems Engineering at Tokyo University of Agriculture and Technology, Japan. His research interests are computer assisted surgery/therapy based on robotics and information science.

#### Takashi Mochizuki

Takashi Mochizuki received the B.Eng. and Ph.D. degrees from Tokyo University of Agriculture and Technology, Japan, in 1976 and 1994, respectively. From 1976 to 2012, he worked for Hitachi Aloka Medical, Ltd. After the company retirement, he established a



startup company to consult the ultrasound technology and has become CEO. Also, he has been a Specially Appointed Professor (NEXT Program, JSPS) of Graduate School of Bio-Applications and Systems Engineering, Tokyo University of Agriculture and Technology. His research interests include therapeutic applications in ultrasound and bioeffects of ultrasound.

# Kohji Masuda

Kohji Masuda received the B.Eng., the M.Eng. and Ph.D. degrees from Osaka University, Japan, in 1991, 1993, and 1996, respectively. From 1996 to 1999, he was an Assistant Professor with Nagoya University, Japan. From 1999 to 2002, he was an



Assistant Professor with Ehime University, Japan. From 2002 to 2003, he was an Postdoctoral Fellow with Universite Grenoble I-Joseph Fourier, France. Since 2003, he has been an Associate Professor with Graduate School of Bio-Applications and Systems Engineering, Tokyo University of Agriculture and Technology, Japan. His research interests include therapeutic applications in ultrasound, medical robotics, and medical image processing.

### Ryo Suzuki

Ryo Suzuki graduated the school of pharmacy, Tokyo university of Pharmacy and Life Sciences, Japan in 1996 and received Ph.D. (Pharmaceutical Sciences) from Graduate School of Pharmaceutical Sciences, Osaka University, Japan in 2001. He is currently an



associated professor, School of Pharma-Sciences, Teikyo University, Japan. His research interest is drug and gene delivery system with liposome and ultrasound.

### Kazuo Maruyama

Kazuo Maruyama graduated School pharmacy, Tokyo University of Pharmacy and life Science, Japan in 1977 and received Master's (Pharmacy) degree from Graduate school of Kumamoto University, Japan, in 1979. He received Ph.D. degrees from Tokyo Uni-



versity of Pharmacy and life Science, Japan, in 1986. He was an Assistant Professor and an Associate Professor with School of Pharmacy, Teikyo University from 1977 to 1988, and from 1994 to 2002, respectively. From 1988 to 1990, he was a Postdoctoral Fellow with University of Tennessee, USA. Since 2002 he has been a Professor with School of Pharmacy, Teikyo University, Japan. His research interests include liposomal drug deliver system, and therapostics which can diagnosis and therapy in combination of nanobubble and ultrasound.

