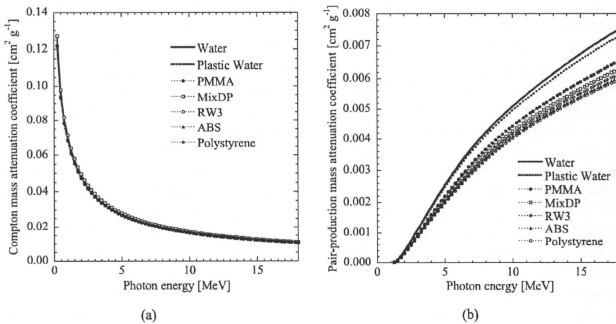


**Table 2.** Depth scaling factors for the open and blocked field and relative electron densities of commercially available solid phantoms. The depth scaling factors are determined from photon energy spectra at 10 cm depth for 10 cm  $\times$  10 cm open and blocked field, respectively.

Phantom	$(\rho_e)_{p/w}$	Nominal photon energy [MV]							
		4		6		10		15	
		Open	*MBF	Open	MBF	Open	MBF	Open	MBF
WT1	0.972	0.971	0.970	0.970	0.969	0.967	0.966	0.964	0.965
457-CTG	0.972	0.971	0.970	0.970	0.969	0.967	0.966	0.964	0.965
RW3	0.966	0.965	0.964	0.964	0.962	0.958	0.957	0.953	0.954
MixDP	1.012	1.010	1.009	1.008	1.006	1.001	1.000	0.996	0.997
WE211	0.973	0.972	0.972	0.971	0.970	0.968	0.967	0.965	0.966
WE211R	0.974	0.974	0.973	0.973	0.972	0.969	0.969	0.966	0.967
Plastic Water	0.969	0.970	0.970	0.970	0.970	0.970	0.970	0.970	0.970
Plastic Water DT	0.963	0.963	0.963	0.962	0.962	0.962	0.962	0.961	0.961
Virtual Water	0.968	0.968	0.967	0.967	0.966	0.963	0.963	0.961	0.961
Polystyrene	0.969	0.967	0.966	0.965	0.963	0.958	0.957	0.953	0.954
PMMA	0.972	0.970	0.970	0.969	0.968	0.965	0.965	0.963	0.963
ABS	0.972	0.970	0.969	0.968	0.966	0.961	0.960	0.956	0.957

\*MBF: MLC blocked field



**Fig. 1.** (a) Compton mass attenuation coefficients and (b) pair-production mass attenuation coefficients for several solid phantoms.<sup>(10)</sup>

mined from the effective mass attenuation coefficient.

#### Depth and field size dependence

Figure 2 shows simulated photon energy spectra in water at various depths (a) on beam central axis for 10 MV photon beam with various size of open field (b) as an example of simulation results. Each spectrum was normalized to total photon fluence of the sampling region.  $E_{av}$  shows average

photon energy determined from the photon spectrum.

As depth increase, photon energy was hardened by attenuation with a medium. The contribution of scattered photon was increased and  $E_{av}$  was decreased with increasing field size. Therefore, the beam hardening effect on the spectrum became small for large field.

Table 3 shows the depth scaling factors of the polystyrene for various depths and field sizes because it shows large

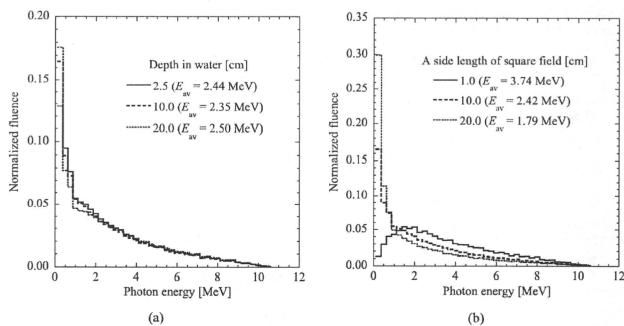


Fig. 2. Photon energy spectra in water for the 10 MV photon beam as a function of (a) depth and (b) field size at 10 cm depth. The each spectrum is normalized to total fluence of the sampling region.  $E_{av}$  shows average photon energies for each photon spectrum.

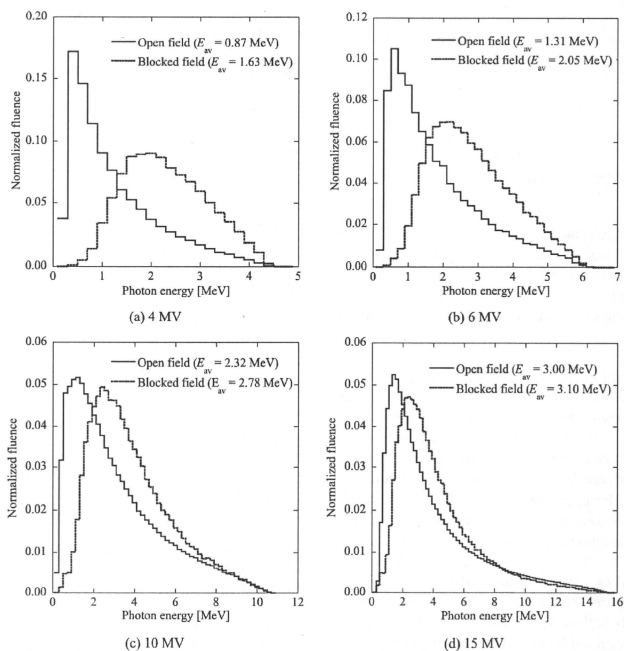
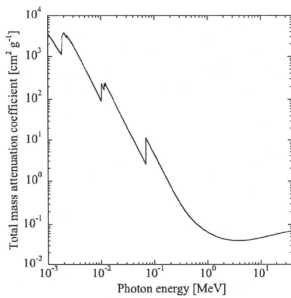


Fig. 3. Photon energy spectra in  $2\text{ cm}\phi$  central area at phantom surface for the open field and blocked field for (a) 4 MV, (b) 6 MV, (c) 10 MV and (d) 15 MV. The each spectrum was normalized to total fluence of the sampling region.

**Table 3.** Depth scaling factors of the polystyrene for various depths and field sizes for 10 MV photon beam.

Depth [cm]	A side length of square field [cm]			
	1.0	5.0	10.0	20.0
2.5	0.956	0.957	0.958	0.958
5.0	0.956	0.957	0.958	0.958
10.0	0.955	0.957	0.958	0.958
15.0	0.955	0.956	0.957	0.958
20.0	0.954	0.955	0.957	0.958



**Fig. 4.** The total mass attenuation coefficient for tungsten.<sup>9)</sup>

photon energy dependency. As mentioned above, the depth scaling factor is strongly depending on proportion of high energy part in the spectrum. The proportion of high energy part was constant even though depth and field size were changed, consequently, depth scaling factors were almost constant.

#### *Influence of beam hardening effect by the MLC*

Photon energy spectra at a phantom surface for open field and blocked field with the same jaw setting, 10 cm × 10 cm, are shown in Fig. 3. For all photon energies, the lower energy photons were removed by the MLC which tungsten has high mass attenuation coefficient for low energy photons as shown in Fig. 4. Furthermore, for 15 MV beam, reduction of high energy photons was also observed because tungsten has minimum mass attenuation coefficient at 4 MeV and the mass attenuation coefficient increases above 4 MeV.

The depth scaling factors for blocked field with 10 cm × 10 cm jaw setting are also tabulated in Table 2. For the Plastic Water, the depth scaling factor for the blocked field was equal to that of open field. On the other hand, the depth scaling factor of the RW3, MixDP, polystyrene, PMMA and ABS for the blocked field were obviously different from the

relative electron densities. For the polystyrene, the difference of  $z_{\text{pl}}$  scaled by the relative electron density becomes 2.0 mm at 10 cm depth. This is because pair-production mass attenuation coefficients of these solid phantoms are different from water. However, the depth scaling factor for the blocked field was almost equal to the depth scaling factor for the open field. The  $z_{\text{pl}}$  determined from both scaling factors agreed within 0.2 mm at 10 cm depth.

Consequently, to minimize the error of the depth scaling for IMRT beams, the depth scaling factor should be obtained from the effective mass attenuation coefficient.

#### *Conclusion*

To clarify the accurate depth scaling method for IMRT beam, the variations of photon energy spectra for various radiation fields and depth scaling factors were determined. The results clarified that the depth scaling factor of Plastic Water were independent on photon energy. On the other hand, the depth scaling factors of the other phantoms were significantly different by photon energy. Therefore depth must be scaled by using presented depth scaling factor corresponding to photon beam energy.

The results also clarified that the ratio of effective mass attenuation coefficients of media to water are unaffected by whether open or blocked field. The presented depth scaling factor can also be used for IMRT beams.

In the future, this study can be used to clarify the effect of spectral variation for fluence scaling.

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